# INDEFINITE EXTRINSIC SYMMETRIC SPACES 

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Abstract. We extend Ferus' characterization of extrinsic symmetric spaces to ambient spaces with indefinite inner products.

## 1. Introduction

Let $V$ be a finite dimensional real vector space with a non-degenerate scalar product (metric) $\langle$,$\rangle . A linear subspace W \subset V$ is called nondegenerate if $\left.\langle\rangle\right|_{W$, remains non-degenerate (which is trivially satisfied if the inner product is positive definite), otherwise $W$ is called degenerate. A smooth submanifold $M \subset V$ is called non-degenerate if so is each of its tangent spaces $T_{x} M \subset V, x \in M$. Then $M$ inherits a semi-Riemannian metric from $V$. Now for any $x \in M$ we consider the reflection at the affine normal space $x+N_{x} M$ where $N_{x} M=T_{x} M^{\perp}$; this is the affine isometry $s_{x}: V \rightarrow V$ with

$$
s_{x}(x)=x,\left.\quad s_{*}\right|_{T_{x} M}=-I,\left.\quad s_{*}\right|_{N_{x} M}=I
$$

A nondegenerate submanifold $M \subset V$ will be called extrinsically symmetric if $s_{x}(M)=M$ for all $x \in M$. Viewed as a semi-Riemannian manifold with the induced metric, $M$ is a symmetric space since $\left.s_{x}\right|_{M}$ is an isometric point reflection for any $x \in M$. Extrinsic symmetric submanifolds are characterized by the property $\nabla \alpha=0$ where $\alpha: S^{2}(T M) \rightarrow N M$ is the second fundamental form and $\nabla \alpha: S^{3}(T M) \rightarrow N M$ its covariant derivative. In fact, every extrinsic symmetric submanifold satisfies this property since each $s_{x}$ preserves $\nabla \alpha_{x}$ but $s_{x}=-I$ on $S^{3}\left(T_{x} M\right)$ (three signs) while $s_{x}=I$ on $N_{x} M$. The converse statement is a theorem of Ferus [4] and Strübing [9].

A rich set of examples is obtained as follows. Let $G$ be a connected Lie group and assume that its Lie algebra $\mathfrak{g}$ is equipped with an $\operatorname{Ad}(G)$-invariant inner product and an orthogonal Cartan decomposition $\mathfrak{g}=\mathfrak{k} \oplus V$, i.e.

$$
\begin{equation*}
[\mathfrak{k}, \mathfrak{k}] \subset \mathfrak{k}, \quad[\mathfrak{k}, V] \subset V, \quad[V, V] \subset \mathfrak{k} . \tag{1}
\end{equation*}
$$

The connected Lie group $K \subset G$ with Lie algebra $\mathfrak{k}$ acts on $V$ by the adjoint action; in fact this is the isotropy representation of the symmetric space $G / K$. Then an orbit $M=\operatorname{Ad}(K) x$ with $x \in V$ is extrinsic symmetric iff

$$
\begin{equation*}
\operatorname{ad}(x)^{3}=\lambda \operatorname{ad}(x) \tag{2}
\end{equation*}
$$

for some $\lambda \neq 0$ (cf. [5], [3]). If $\lambda=-1$, the extrinsic symmetry at $x$ is $s_{x}=$ $\exp (\pi \operatorname{ad}(x))$. For $\lambda<0$, this is just a normalization of $x$; for $\lambda>0$ we pass to the dual Lie algebra $\mathfrak{g}^{*}=\mathfrak{k} \oplus i V \subset \mathfrak{g} \otimes \mathbb{C}$ where $i=\sqrt{-1}$. These extrinsic symmetric spaces will be called of Ferus type (see below). Most Riemannian symmetric spaces (the so called symmetric $R$-spaces, including Grassmannians, conjugacy classes of real and complex structures, hermitian symmetric spaces and the Lie groups $S O_{n}$,

[^0]$U_{n}, S p_{n}$ ) can be isometrically embedded as extrinsic symmetric submanifolds of Ferus type.

In general, if $M \subset V$ is extrinsic symmetric, we let

$$
\begin{equation*}
\hat{K}=\left\langle s_{x} ; x \in M\right\rangle \subset O(V) \tag{3}
\end{equation*}
$$

be the group generated by all reflection $s_{x}$. The affine map $s_{x}$ will be called extrinsic symmetry at $x$, and the group $\hat{K}$ is the symmetry group of $M$. We denote its connected component $K$ (the transvection group) and its Lie algebra by $\mathfrak{k}$. We call $M \subset V$ full if it does not lie in a proper affine subspace, and indecomposable if there is no nontrivial orthogonal splitting $V=V_{1} \oplus V_{2}$ and $M=M_{1} \times M_{2}$ with $M_{i} \subset V_{i}$.

Dirk Ferus has given the following characterization of extrinsic symmetric spaces in the case where the inner product is positive definite:
Theorem: (Ferus [4],[5]) Let $M \subset V$ (1) full, (2) indecomposable, (3) extrinsic symmetric. Then we have:
(A) The vector space $\mathfrak{g}:=\mathfrak{k} \oplus V$ carries the structure of a Lie algebra with Cartan decomposition, and the action of $K$ on $V$ is the the adjoint action of $K$ restricted to $V \subset \mathfrak{g}$.
(B) $M$ is the $K$-orbit of an element $x \in V$ with $\operatorname{ad}(x)^{3}=\lambda \operatorname{ad}(x)$ for some $\lambda \neq 0$ (i.e. it is of Ferus type).

It is the aim of the present paper to generalize this theorem to the case where the inner product is indefinite. Let us briefly discuss the assumption of Ferus' theorem under this view point. On the one hand, the fullness assumption (1) seems too strong; one would like to discuss the case where $M$ is contained in a proper subspace $W \subset V$ where the inner product is degenerate (otherwise we could just pass to $W$ in place of $V)$. However, projecting $M \subset W$ onto the quotient vector space $W /\left.\operatorname{ker} s\right|_{W}$ where the induced inner product is nondegenerate, we restore the assumption of the theorem, see Section 2. On the other hand, the indecomposability assumption (2) is much weaker in the indefinite case since it allows for nontrivial (degenerate) $K$-invariant affine subspaces of $V$ which even may be contained in $M$. However, we are still able to prove essential parts of Ferus' theorem. The result involves the shape operator (Weingarten map) $S_{H}(v)=-\partial_{v} H$ for the mean curvature vector $H=\operatorname{tr} \alpha$ where $\alpha$ is the second fundamental form of $M$, i.e. $\alpha(v, w)=\left(\partial_{v} w\right)^{N}$.
Theorem A. Let $M \subset V$ be a (possibly immersed) nondegenerate submanifold which is (1) full, (2) indecomposable, (3) extrinsic symmetric. Then the vector space $\mathfrak{g}:=\mathfrak{k} \oplus V$ carries the structure of a Lie algebra with Cartan decomposition, and the linear part of the (affine) action of $K$ on $V$ is the the adjoint action of $K$ restricted to $V \subset \mathfrak{g}$.

Theorem B. Under the above assumptions, $M$ is of Ferus type unless $S_{H}^{2}=0$.
The present work is based on the 2005 thesis [7] of the first named author. Previously, Naitoh [8] had proved Theorem B by a different method under the additional assumption that there is an umbilic normal vector field. Recently, I. Kath [6] has given a full description of indefinite extrinsic symmetric spaces.
It is our pleasure to express our warmest thanks to Ines Kath for many suggestions and helpful conversations. We also would like to thank the referee for his very helpful suggestions.

## 2. The fullness assumption

Let us assume that $M \subset V$ is an extrinsic symmetric space lying in some proper linear subspace $W \subset V$. If the inner product is nondegenerate on $W$, we may
replace $W$ by $V$. The interesting case is when $\left.\langle\rangle\right|_{W$,$} is degenerate, but still M$ is nondegenerate. Let

$$
N=\left.\operatorname{ker}\langle,\rangle\right|_{W}=\left\{n \in W ;\langle n, w\rangle=0 \forall_{w \in W}\right\}
$$

Then $\bar{W}=W / N$ inherits a nondegenerate inner product from $\left.\langle\rangle\right|_{W$,$} . Since any$ tangent space $T_{x} M$ is nondegenerate, it does not intersect $N$, and thus the canonical projection $\pi: W \rightarrow \bar{W}$ is an immersion on $M$. Moreover, $\pi$ is an isometry, hence it conjugates the reflection at the normal spaces of $M \subset W$ and $\pi(M) \subset \bar{W}$ which shows that $\pi: M \rightarrow \bar{W}$ is extrinsic symmetric. Hence we have proved:

Theorem 2.1. Let $W$ be a real vector space with a possibly degenerate inner product $s=\langle$,$\rangle and M \subset W$ a (possibly immersed) extrinsic symmetric submanifold. Let $N=\operatorname{ker} s$ and let $\pi: W \rightarrow \bar{W}=W / N$ be the canonical projection. Then the vector space $\bar{W}$ inherits a nondegenerate inner product which makes $\pi$ isometric, and $\left.\pi\right|_{M}: M \rightarrow \bar{W}$ is an extrinsic symmetric immersion.

A simple example is $W=\mathbb{R}^{2}$ with the inner product $\langle x, y\rangle=x_{1} y_{1}$ where we let $M$ be the graph of a parabola, $M=\left\{\left(u, t u^{2}\right) ; u \in \mathbb{R}\right\}$ for arbitrary $t \in \mathbb{R}$. For any $x=\left(u, t u^{2}\right) \in M$ we have $T_{x} M=\mathbb{R}(1,2 t u)$ while $N_{x} M$ is always the vertical line $N=\mathbb{R}(0,1)$. The reflection $s_{x}$ fixes $x$ while its differential $\left(s_{x}\right)_{*}$ fixes the vector $e_{2}=(0,1)$ and maps $(1,2 t u)$ to $-(1,2 t u)$. Thus for any $y=\left(v, t v^{2}\right) \in M$ we have $s_{x}\left(v, t v^{2}\right)=\left(w, t w^{2}\right) \in M$ with $w=2 u-v$, hence $M$ is extrisic symmetric. The projection $\pi:(u, v) \mapsto u$ maps $M$ onto the real line (the $x_{1}$-axis) with its trivial extrinsic symmetric structure.


Remark. More generally, suppose we have a nondegenerate inner product space $\bar{W}$ containing an (immersed) extrinsic symmetric submanifold $\bar{M} \subset \bar{W}$. Let $W=$ $\bar{W} \oplus N$ with zero inner product on $N$, and let $\pi: W \rightarrow \bar{W}$ be the projection onto the first factor. We may ask for all (immersed) extrinsic symmetric $M \subset W$ with $\pi(M)=\bar{M}$. As the above example suggests, this is a nontrivial problem, cf. [6].

## 3. Constructing the Lie bracket

From now on we assume fullness. The idea for the proof of Theorem A is to adapt a method of [3] which gives an alternative proof of Ferus' Theorem for positive definite inner products. Let $M \subset V$ be full and extrinsic symmetric and $K$ its symmetry group as above. Fix some $x \in M$. Let $\mathfrak{k}=\mathfrak{k}_{+}+\mathfrak{k}_{-}$be the Cartan decomposition of $\mathfrak{k}$ corresponding to the extrinsic symmetry $s_{x}$, i.e. conjugation with $s_{x}$ on $\mathfrak{k}$ fixes $\mathfrak{k}_{+}$and anti-fixes $\mathfrak{k}_{-}$. The linear map sending $A \in \mathfrak{k}_{-}$onto $A x \in T_{x} M$ is an $K$-equivariant isomorphism between $\mathfrak{k}_{-}$and $T_{x} M$; its inverse will be called $v \mapsto t_{v}: T_{x} M \rightarrow \mathfrak{k}_{-}$. In fact, the map $t_{v}$ is the infinitesimal transvection in $v$-direction, i.e. the affine isometries $k_{v}(t)=\exp t t_{v} \in K$ form the one-parameter group of transvections along the geodesic $\gamma_{v}$ whose differential is the parallel transport along $\gamma$ in both the tangent and normal bundles. Obviously, the linearized
action of $\mathfrak{k}_{+}$(being the Lie algebra of the isotropy group $K_{+}=K_{x}$ ) preserves the splitting $V=V_{+}+V_{-}$where

$$
\begin{equation*}
V_{+}=N_{x} M, \quad V_{-}=T_{x} M \tag{4}
\end{equation*}
$$

On the other hand, the linearized action of $\mathfrak{k}_{-}$reverses this splitting since by differentiating a Levi-Civita-parallel tangent (resp. normal) vector field along $\gamma_{v}$ we obtain a normal (resp. tangent) vector field. More precisely, for any $v, w \in V_{-}$and $\xi \in V_{+}$we have

$$
\begin{equation*}
t_{v} w=\alpha(v, w), \quad t_{v} \xi=-S_{\xi}(v) \tag{5}
\end{equation*}
$$

where $S_{\xi}: V_{-} \rightarrow V_{-}$denotes the shape operator or Weingarten map defined by

$$
\begin{equation*}
S_{\xi}(v)=-\left(\partial_{v} \xi\right)^{T}, \quad\left\langle S_{\xi}(v), w\right\rangle=\langle\alpha(v, w), \xi\rangle . \tag{6}
\end{equation*}
$$

Since the linear maps $t: \mathfrak{p}_{-} \rightarrow \mathfrak{k}_{-}, v \mapsto t_{v}$ and $S: V_{+} \rightarrow S\left(V_{-}\right), \xi \mapsto S_{\xi}$ are equivariant with respect to the action of $\mathfrak{k}_{+}$, we have for all $A \in \mathfrak{k}_{+}$:

$$
\begin{equation*}
\left[A, t_{v}\right]=t_{A_{*} v}, \quad\left[A_{*}, S_{\xi}\right]=S_{A_{*} \xi} \tag{7}
\end{equation*}
$$

First we construct a certain $\operatorname{Ad}(K)$-invariant inner product on $\mathfrak{k}$. Using the $K_{+}-$ equivariant isomorphism $T: V_{-} \rightarrow \mathfrak{k}_{-}$we may transplant the inner product on $V_{-}$ to $\mathfrak{k}_{-}$:

$$
\begin{equation*}
\left\langle t_{v}, t_{w}\right\rangle_{\mathfrak{k}_{-}}=\langle v, w\rangle \tag{8}
\end{equation*}
$$

This is extended to all of $\mathfrak{k}$ by declaring $\mathfrak{k}_{+} \perp \mathfrak{k}_{-}$and defining the metric on $\mathfrak{k}_{+}$as follows:

$$
\begin{equation*}
\left\langle A,\left[t_{v}, t_{w}\right]\right\rangle_{\mathfrak{k}_{+}}=\left\langle\left[A, t_{v}\right], t_{w}\right\rangle_{\mathfrak{k}_{-}}=\langle A v, w\rangle \tag{9}
\end{equation*}
$$

for any $v, w \in T_{x} M$. In fact, according to the subsequent lemma we have $\left[\mathfrak{k}_{-}, \mathfrak{k}_{-}\right]=$ $\mathfrak{k}_{+}$, hence $\left[t_{v}, t_{w}\right]$ is a generic element of $\mathfrak{k}_{+}$. We must show that the metric is well defined by (9), in other words that $\langle A v, w\rangle$ only depends on $B:=\left[t_{v}, t_{w}\right]$. We will use the general formula

$$
\begin{equation*}
\left[t_{a}, t_{b}\right] c=-R(a, b) c \tag{10}
\end{equation*}
$$

for the curvature tensor $R$ of the symmetric space $M$. Putting $A=\left[t_{a}, t_{b}\right]$ for some $a, b \in T_{x} M$ we find $\langle A v, w\rangle=\left\langle\left[t_{a}, t_{b}\right] v, w\right\rangle=-\langle R(a, b) v, w\rangle=-\langle R(v, w) a, b\rangle=$ $\left\langle\left[t_{v}, t_{w}\right] a, b\right\rangle=\langle B a, b\rangle$. Clearly, the metric is $\mathfrak{k}_{+}$-invariant, and from (9) it is also $\mathfrak{k}_{-}$-invariant since $\langle A v, w\rangle=\left\langle\left[A, t_{v}\right], t_{w}\right\rangle$.

It remains to show that this inner product is non-degenerate. Since the $\mathfrak{k}_{-}$-part is nondegenerate (being a copy of the $V_{-}$-part), it is enough to prove nondegeneracy for the $\mathfrak{k}_{+}$-part. If this is degenerate, there exists a nonzero $A_{o} \in \mathfrak{k}_{+}$with $\left\langle A_{o} v, w\right\rangle=$ 0 for all $v, w \in T_{x} M$, cf. (9), hence $A_{o} v=0$ for all $v$ which contradicts to the effectivity of the isotropy action.

Lemma 3.1. (cf. [2]) $\left[\mathfrak{k}_{-}, \mathfrak{k}_{-}\right]=\mathfrak{k}_{+}$.
Proof. This is only due to the fact that $K^{o}$ is generated by the transvections. In fact, let $\mathfrak{k}_{1}=\mathfrak{k}_{-}+\left[\mathfrak{k}_{-}, \mathfrak{k}_{-}\right]$. From the Cartan relations $\left[\mathfrak{k}_{-}, \mathfrak{k}_{+}\right] \subset \mathfrak{k}_{-}$and $\left[\mathfrak{k}_{-}, \mathfrak{k}_{-}\right] \subset \mathfrak{k}_{+}$ we have $\left[\mathfrak{k}_{-},\left[\mathfrak{k}_{-}, \mathfrak{k}_{-}\right]\right] \subset\left[\mathfrak{k}_{-}, \mathfrak{k}_{+}\right] \subset \mathfrak{k}_{-}$(Lie triple property) and therefore $\mathfrak{k}_{1} \subset \mathfrak{k}$ is a Lie subalgebra. Let $K_{1} \subset K$ be the corresponding connected Lie subgroup. We need to show that all transvections $s_{y} s_{z}$ for any two $y, z \in M$ belong to $K_{1}$. It suffices to show this when $y$ and $z$ are connected by a geodesic $\gamma$ : In general there is only a geodesic polygon connecting $y$ and $z$, but labeling its vertices $y=y_{0}$, $y_{1}, \ldots, y_{k}=z$, we have $s_{y} s_{z}=s_{y_{0}} s_{y_{1}} s_{y_{1}} s_{y_{2}} \ldots s_{y_{k-1}} s_{y_{k}}$ and hence it suffices to show $s_{y_{i-1}} s_{y_{i}} \in K_{1}$.

Thus we assume that $y$ and $z$ lie on a common geodesic $\gamma$ with (say) $\gamma(0)=z$ and $\gamma(1 / 2)=y$. We want to show $s_{y} s_{z} \in K_{1}$. If $z=x$, this is obvious since then $s_{y} s_{x}=\exp t_{v} \subset \exp \mathfrak{k}_{-}$where $v=\gamma^{\prime}(0)$. For an arbitrary geodesic segment $\gamma:[0,1] \rightarrow M$ we put

$$
\tau_{\gamma}=s_{\gamma(1 / 2)} s_{\gamma(0)}
$$

This is the transvection along $\gamma$ sending $\gamma(0)$ to $\gamma(1)$. Likewise for a geodesic polygon $p=\gamma_{1} * \cdots * \gamma_{k}$ (concatenation of $k$ geodesic segments) we put $\tau_{p}=$ $\tau_{\gamma_{k}} \ldots \tau_{\gamma_{1}}$. By induction over $k$ we claim $\tau_{p} \in K_{1}$ for any geodesic polygon $p$ starting at our base point $x$. This is clear for $k=1$. In the general case we let $p=p^{\prime} * \gamma$ where $\gamma=\gamma_{k}$ and $p^{\prime}=\gamma_{1} * \cdots * \gamma_{k-1}$. By induction hypothesis $\tau_{p^{\prime}} \in K_{1}$. Further, the geodesic $\beta=\tau_{p^{\prime}}^{-1} \circ \gamma$ starts at $x$, and hence $\tau_{\beta} \in K_{1}$ whence $\tau_{\gamma}=\tau_{p^{\prime}}^{-1} \tau_{\beta} \tau_{p^{\prime}} \in K_{1}$.

Now let $y, z \in M$ be as above. We have to make sure that $s_{y} s_{z}=\tau_{\gamma}$ is contained in $K_{1}$. Join $x$ to $z$ by a geodesic polygon $p$. Then $\tau_{p} \in K_{1}$, and the geodesic $\beta=\tau_{p}^{-1} \circ \gamma$ starts at $x$. Thus $\tau_{\beta} \in K_{1}$ whence $\tau_{\gamma}=\tau_{p} \tau_{\beta} \tau_{p}^{-1} \in K_{1}$.

Next we define a skew symmetric product [, ] on

$$
\begin{equation*}
\mathfrak{g}=\mathfrak{k}+V \tag{11}
\end{equation*}
$$

extending the Lie product on $\mathfrak{k}$ as follows. For any $A \in \mathfrak{k}$ and $v, w \in V$ we define $[A, v] \in V$ and $[v, w] \in \mathfrak{k}$ by

$$
\begin{equation*}
[A, v]=A_{*} v, \quad\langle[v, w], A\rangle_{\mathfrak{k}}=\langle A v, w\rangle . \tag{12}
\end{equation*}
$$

Note that

$$
\begin{equation*}
[v, w]=\left[t_{v}, t_{w}\right]=R(v, w), \tag{13}
\end{equation*}
$$

since $\left\langle A,\left[t_{v}, t_{w}\right]\right\rangle=\left\langle\left[A, t_{v}\right], t_{w}\right\rangle=\left\langle t_{A v}, t_{w}\right\rangle=\langle A v, w\rangle$.
By [3], p. 518-519, we have defined a Lie algebra stucture on $\mathfrak{g}$ (the main work amounts to proving the Jacobi identity for all $u, v, w \in V \subset \mathfrak{g}$ ) with Cartan decomposition (11), and the direct sum metric on $\mathfrak{g}=\mathfrak{k}+V$ is ad $(\mathfrak{g})$-invariant. Further,

$$
\begin{equation*}
\mathfrak{g}_{+}:=\mathfrak{k}_{+} \oplus V_{+} \tag{14}
\end{equation*}
$$

is a subalgebra since $\left\langle\left[V_{+}, V_{+}\right], \mathfrak{k}_{-}\right\rangle_{\mathfrak{g}}=\left\langle\left[\mathfrak{k}_{-}, V_{+}\right], V_{+}\right\rangle_{\mathfrak{g}} \subset\left\langle V_{-}, V_{+}\right\rangle=0$. Putting

$$
\begin{equation*}
\mathfrak{g}_{-}=\mathfrak{k}_{-} \oplus V_{-}, \tag{15}
\end{equation*}
$$

we have a second Cartan decomposition,

$$
\begin{equation*}
\mathfrak{g}=\mathfrak{g}_{+}+\mathfrak{g}_{-} \tag{16}
\end{equation*}
$$

which is compatible to the first one. This finishes the proof of Theorem A.

## 4. Linearity of the action

Theorem 4.1. Let $M \subset V$ be full, indecomposable and extrinsic symmetric. Suppose that all affine normal spaces $x+N_{x} M, x \in M$, have a nonempty intersection. Then $M$ is of Ferus type.

Proof. Up to translations we may assume that the common intersection of the affine normal spaces contains the origin 0 . Thus $x \in N_{x} M$ for all $x \in M$, and the extrinsic symmetries $s_{x}$ are linear (fixing 0 ). Hence $\hat{K}=\left\langle s_{x} ; x \in M\right\rangle$ is a subgroup of the orthogonal group $O(V)$. Now we fix $x \in M$ and let $V_{-}=T_{x} M$ and $V_{+}=N_{x} M$. Then for all $v, w \in V_{-}$we have

$$
\begin{aligned}
\operatorname{ad}(x) t_{v} & =-t_{v} x=-v \\
\left\langle t_{w}, \operatorname{ad}(x) v\right\rangle & =\left\langle t_{w} x, v\right\rangle=\langle w, v\rangle=\left\langle t_{w}, t_{v}\right\rangle .
\end{aligned}
$$

Further, $\mathfrak{k}_{+}$is the stabilizer subalgebra for $x$, hence for all $A \in \mathfrak{k}_{+}$and $\xi \in V_{+}$we have $[A, x]=0$ and $\langle A,[x, \xi]\rangle=\langle[A, x], \xi\rangle=0$. Thus

$$
\begin{equation*}
\operatorname{ad}(x) v=t_{v}, \quad \operatorname{ad}(x) t_{v}=-v,\left.\quad \operatorname{ad}(x)\right|_{\mathfrak{g}_{+}}=0 \tag{17}
\end{equation*}
$$

showing $M=\operatorname{Ad}(K) x$ with $\operatorname{ad}(x)^{3}=-\operatorname{ad}(x)$ which shows that $M$ has Ferus type.

Remark. The situation looks very similar to the case of positive definite inner product. However note that $x$ could be a light vector, i.e. $\langle x, x\rangle=0$. Otherwise, if $\langle x, x\rangle=s \neq 0$, then $M$ is contained in the "sphere" $S_{s}=\{x \in V ;\langle x, x\rangle=s\}$, but different from the positive definite case, it need not to be minimal inside $S_{s}$.

## 5. The Killing form

We have constructed a metric Lie algebra $\mathfrak{g}=\mathfrak{k}+V$ equipped with two commuting involutions $\sigma, \tau$ leading to the orthogonal decomposition

$$
\begin{equation*}
\mathfrak{g}=\mathfrak{k}_{+}+\mathfrak{k}_{-}+V_{+}+V_{-} . \tag{18}
\end{equation*}
$$

A very important tool is the Killing form $B$ on the Lie algebra $\mathfrak{g}$ which turns out to be closely related to the second fundamental form $\alpha$ of $M \subset V$. Recall that for any $X, Y \in \mathfrak{g}$ we have

$$
\begin{equation*}
B(X, Y)=\operatorname{tr}_{\mathfrak{g}} \operatorname{ad}(X) \operatorname{ad}(Y)=\sum_{i} \epsilon_{i}\left\langle\operatorname{ad}(X) \operatorname{ad}(Y) E_{i}, E_{i}\right\rangle_{\mathfrak{g}} \tag{19}
\end{equation*}
$$

where $\left(E_{i}\right)$ is an orthonormal basis of $\mathfrak{g}$, i.e. $\left\langle E_{i}, E_{j}\right\rangle=\epsilon_{i} \delta_{i j}$ with $\epsilon_{i}= \pm 1$. We will choose this basis adapted to the orthogonal splitting (18).
Lemma 5.1. For any $v, w \in V_{-}$and $\xi, \eta \in V_{+}$we have

$$
\begin{aligned}
B(v, w) & =2 \operatorname{tr}_{V_{-}}(\operatorname{ad}(v) \operatorname{ad}(w))+2 \operatorname{tr}_{\mathfrak{k}_{-}}(\operatorname{ad}(v) \operatorname{ad}(w)), \\
B\left(t_{v}, t_{w}\right) & =2 \operatorname{tr}_{\mathfrak{k}_{-}}\left(\operatorname{ad}\left(t_{v}\right) \operatorname{ad}\left(t_{w}\right)\right)+2 \operatorname{tr}_{V_{-}}\left(\operatorname{ad}\left(t_{v}\right) \operatorname{ad}\left(t_{w}\right)\right), \\
B(\xi, \eta) & =2 \operatorname{tr}_{\mathfrak{k}_{-}}(\operatorname{ad}(\xi) \operatorname{ad}(\eta))+2 \operatorname{tr}_{\mathfrak{k}_{+}}(\operatorname{ad}(\xi) \operatorname{ad}(\eta)) .
\end{aligned}
$$

Proof. Due to the Cartan relations

$$
\begin{equation*}
[V, V] \subset \mathfrak{k}, \quad\left[\mathfrak{g}_{-} \mathfrak{g}_{-}\right] \subset \mathfrak{g}_{+}, \tag{20}
\end{equation*}
$$

the linear map $\operatorname{ad}(v)$ is a skew symmetric transformation which maps $V_{ \pm}$to $\mathfrak{k}_{\mp}$ and vice versa. Likewise $\operatorname{ad}\left(t_{v}\right)$ interchanges the two subspaces in each of the pairs $\left(V_{+}, V_{-}\right)$and $\left(\mathfrak{k}_{+}, \mathfrak{k}_{-}\right)$, and the same holds for $\operatorname{ad}(\xi)$ and the pairs $\left(\mathfrak{k}_{-}, V_{-}\right)$ and $\left(\mathfrak{k}_{+}, V_{+}\right)$. This shows that the partial traces for $B$ are the same on the two components of each pair: E.g. on $V_{+}+\mathfrak{k}_{-}$we have $\operatorname{ad}(v)=\left({ }_{A} A^{\prime}\right)$ where $A=$ $\left.\operatorname{ad}(v)\right|_{V_{+}}$and $A^{\prime}=\left.\operatorname{ad}(v)\right|_{\mathfrak{e}_{-}}$, but since $\operatorname{ad}(v)^{*}=-\operatorname{ad}(v)$, we have $A^{\prime}=-A^{*}$. Thus

$$
\begin{aligned}
\operatorname{tr}_{V_{+}+\mathfrak{k}_{-}}(\operatorname{ad}(v) \operatorname{ad}(w)) & =\operatorname{tr}_{V_{+}+\mathfrak{k}_{-}}\left(A-A^{*}\right)\left(B_{B}^{-B^{*}}\right) \\
& =\operatorname{tr}_{V_{+}+\mathfrak{k}_{-}}\left(-A^{*} B-A B^{*}\right) \\
& =-\operatorname{tr}\left(A^{*} B\right)-\operatorname{tr}\left(A B^{*}\right),
\end{aligned}
$$

and the latter two terms are equal. Using a similar argument on $V_{-}+\mathfrak{k}_{+}$we obtain

$$
\begin{aligned}
& \operatorname{tr}_{V_{+}+\mathfrak{k}_{-}}(\operatorname{ad}(v) \operatorname{ad}(w))=2 \operatorname{tr}_{\mathfrak{k}_{-}}(\operatorname{ad}(v) \operatorname{ad}(w)), \\
& \operatorname{tr}_{V_{-}+\mathfrak{k}_{+}}(\operatorname{ad}(v) \operatorname{ad}(w))=2 \operatorname{tr}_{V_{-}}(\operatorname{ad}(v) \operatorname{ad}(w)
\end{aligned}
$$

which shows the equality for $B(v, w)=\operatorname{trg}_{\mathfrak{g}}(\operatorname{ad}(v) \operatorname{ad}(w))$. The other two equations follow quite similarly.

Lemma 5.2. For any $v, w \in V_{-}$we have

$$
\begin{equation*}
B(v, w)=-2\langle\alpha(v, w), H\rangle=B\left(t_{v}, t_{w}\right) \tag{21}
\end{equation*}
$$

where $H=\operatorname{tr}_{V_{-}} \alpha$.
Proof. Choosing an orthonormal basis $\left(e_{1}, \ldots, e_{m}\right)$ of $V_{-}$, we get

$$
\begin{aligned}
B(v, w) & =2 \sum \epsilon_{i}\left\langle\operatorname{ad}(v) \operatorname{ad}(w) e_{i}, e_{i}\right\rangle+2 \sum \epsilon_{i}\left\langle\operatorname{ad}(v) \operatorname{ad}(w) t_{e_{i}}, t_{e_{i}}\right\rangle \\
& =-2 \sum \epsilon_{i}\left\langle\left[w, e_{i}\right],\left[v, e_{i}\right]\right\rangle-2 \sum \epsilon_{i}\left\langle t_{e_{i}} w, t_{e_{i}} v\right\rangle \\
& =2 \sum \epsilon_{i}\left\{\left\langle R\left(w, e_{i}\right) v, e_{i}\right\rangle-\left\langle\alpha\left(e_{i}, w\right), \alpha\left(e_{i}, v\right)\right\rangle\right\}
\end{aligned}
$$

$$
\begin{aligned}
& \stackrel{(G)}{=}-2 \sum \epsilon_{i}\left\langle\alpha\left(e_{i}, e_{i}\right), \alpha(v, w)\right\rangle=-2\langle H, \alpha(v, w)\rangle, \\
B\left(t_{v}, t_{w}\right) & =2 \sum \epsilon_{i}\left\langle\operatorname{ad}\left(t_{v}\right) \operatorname{ad}\left(t_{w}\right) t_{e_{i}}, t_{e_{i}}\right\rangle+2 \sum \epsilon_{i}\left\langle\operatorname{ad}\left(t_{v}\right) \operatorname{ad}\left(t_{w}\right) e_{i}, e_{i}\right\rangle \\
& =-2 \sum \epsilon_{i}\left\langle\left[t_{w}, t_{e_{i}}\right],\left[t_{v}, t_{e_{i}}\right]\right\rangle-2 \sum \epsilon_{i}\left\langle t_{w} e_{i}, t_{v} e_{i}\right\rangle \\
& =2 \sum \epsilon_{i}\left\{\left\langle R\left(w, e_{i}\right) v, e_{i}\right\rangle-\left\langle\alpha\left(e_{i}, w\right), \alpha\left(e_{i}, v\right)\right\rangle\right\} \\
& \stackrel{(G)}{=}-2 \sum \epsilon_{i}\left\langle\alpha\left(e_{i}, e_{i}\right), \alpha(v, w)\right\rangle=-2\langle H, \alpha(v, w)\rangle,
\end{aligned}
$$

where $(G)$ refers to the Gauss equations for $M \subset V$,

$$
\begin{equation*}
\langle R(a, b) c, d\rangle=\langle\alpha(a, d), \alpha(b, c)\rangle-\langle\alpha(b, d), \alpha(a, c)\rangle . \tag{G}
\end{equation*}
$$

This finishes the proof.
Lemma 5.3. Let $\xi, \eta \in V_{+}$with $\xi=\alpha(v, w)=t_{v} w$. Then

$$
\begin{equation*}
B(\xi, \eta)=-2\left\langle\alpha\left(S_{\eta} v, w\right), H\right\rangle \tag{22}
\end{equation*}
$$

Proof. $B(\xi, \eta)=B\left(t_{v} w, \eta\right)=-B\left(w, t_{v} \eta\right)=B\left(w, S_{\eta} v\right) \stackrel{(21)}{=}-2\left\langle\alpha\left(w, S_{\eta} v\right), H\right\rangle$.
Lemma 5.4. Let $\xi, \eta \in V_{+}$with $\left[\mathfrak{k}_{+}, \eta\right]=0$. Then

$$
\begin{equation*}
B(\xi, \eta)=-2 \operatorname{tr}_{V_{-}}\left(S_{\xi} S_{\eta}\right) \tag{23}
\end{equation*}
$$

Proof. From Lemma 5.1 we obtain

$$
B(\xi, \eta)=-2 \sum \epsilon_{i}\left\langle\left[\xi, t_{e_{i}}\right],\left[\eta, t_{e_{i}}\right]\right\rangle-2 \sum \epsilon_{j}\left\langle\left[\xi, A_{j}\right],\left[\eta, A_{j}\right]\right\rangle
$$

where $\left(A_{j}\right)$ denotes an orthonormal basis of $\mathfrak{k}_{+}$. But the second term vanishes by the assumption $\left[\eta, \mathfrak{k}_{+}\right]=0$, hence the claim follows using (5).

## 6. The shape operators

Let $M \subset V$ be extrinsic symmetric. As before, we fix some $x \in M$ and let $V_{-}=T_{x} M$ and $V_{+}=N_{x} M$. For all normal vectors $\xi \in V_{+}$we consider the shape operators $S_{\xi}$ which are self adjoint endomorphisms of $V_{-}$. Their commutators are obtained from the Ricci equation

$$
\begin{equation*}
\left\langle R^{N}(v, w) \eta, \xi\right\rangle=\left\langle\left[S_{\xi}, S_{\eta}\right] v, w\right\rangle \tag{R}
\end{equation*}
$$

for any $v, w \in V_{-}$and $\xi, \eta \in V_{+}$where $R^{N}$ denotes the curvature tensor of the normal bundle. If $\eta$ can be extended to a normal vector field which is parallel, the left hand side of $(R)$ vanishes, hence the corresponding shape operator $S_{\eta}$ commutes with any $S_{\xi}$.

Most important among the normal vectors is the mean curvature vector

$$
\begin{equation*}
H=\operatorname{tr} \alpha=\sum \epsilon_{i} \alpha\left(e_{i}, e_{i}\right) \tag{24}
\end{equation*}
$$

where $e_{1}, \ldots, e_{m}$ is an orthonormal basis of $V_{-}$, i.e. $\left\langle e_{i}, e_{j}\right\rangle=\epsilon_{i} \delta_{i j}$ with $\epsilon_{i}= \pm 1$. Since $\alpha$ is parallel, so are $H$ and $S_{H}$, and hence $S_{H}$ commutes with any $S_{\xi}$ (Ricci equation).

Yet there are other parallel normal fields closely related to $H$, e.g.

$$
\begin{equation*}
H_{1}:=\operatorname{tr} \alpha\left(S_{H} ., .\right)=\sum \epsilon_{i} \alpha\left(S_{H} e_{i}, e_{i}\right) . \tag{25}
\end{equation*}
$$

## Lemma 6.1.

$$
\begin{equation*}
S_{H_{1}}=\left(S_{H}\right)^{2} . \tag{26}
\end{equation*}
$$

Proof. We apply the lemmas 5.3 and 5.4 for $\xi=\alpha(v, w)$ and $\eta=H$. The two expressions for $-B(\xi, H)$ are

$$
\begin{aligned}
-B(\xi, H) & =2\left\langle\alpha\left(S_{H} v, w\right), H\right\rangle=2\left\langle S_{H} S_{H} v, w\right\rangle \\
-B(\xi, H) & =2 \sum \epsilon_{i}\left\langle S_{\xi} e_{i}, S_{H} e_{i}\right\rangle \\
& =2 \sum \epsilon_{i}\left\langle\alpha\left(e_{i}, S_{H} e_{i}\right), \xi\right\rangle \\
& =2\left\langle H_{1}, \xi\right\rangle \\
& =2\left\langle H_{1}, \alpha(v, w)\right\rangle=2\left\langle S_{H_{1}} v, w\right\rangle .
\end{aligned}
$$

In the case where the metric on $V$ is positive definite, $S_{H}$ is diagonalizable with constant real eigenvalues $\lambda$, and $T M$ splits into the mutually orthogonal eigendistributions $E_{\lambda}$. Since $S_{H}$ is parallel, so are the $E_{\lambda}$, and any shape operator $S_{\xi}$ preserves $E_{\lambda}$. Hence by indecomposability (Moore Lemma) there can be only one eigenvalue $\lambda$. This must be nonzero since $M$ cannot be minimal (cf. [1]). ${ }^{1}$ Now the map $M \ni x \mapsto x+\lambda^{-1} H(x) \in V$ is constant which shows that $M$ is contained in a sphere $S \subset V$ and $M$ is of Ferus type.

In the indefinite case we arrive at the same conclusion when $S_{H}$ is diagonalizable with real eigenvalues. But this cannot be concluded from the symmetry of $S_{H}$ anymore. Therefore, as a first step, we replace the eigenspace decomposition by a coarser one: For each eigenvalue $\lambda \in \mathbb{C}$ of $S_{H}$ we put

$$
E_{\lambda}=\operatorname{ker}\left(S_{H}-\lambda I\right)^{k}
$$

for a sufficiently large integer $k$. These generalized eigenspaces form the Jordan decomposition of $V^{c}=V \otimes \mathbb{C}$ which by self adjointness of $T=\left(S_{H}-\lambda I\right)^{k}$ is orthogonal with respect to the complexified inner product. ${ }^{2}$ A real decomposition is obtained by combining conjugate pairs of eigenvalues $\lambda$ and $\bar{\lambda}$. The subspaces $E_{\lambda}+E_{\bar{\lambda}}$ form a real parallel decomposition of $T M$ which again is $S_{\xi}$-invariant for any normal vector $\xi$. Using the lemma of Moore (cf. [1]), we may conclude from the indecomposability that there is just one conjugate pair $\lambda, \bar{\lambda}$ of eigenvalues of $S_{H}$ (which of course might be equal).

But we can do better. We consider the space $P$ of all parallel real normal fields on $M$. Since all shape operators $S_{\eta}$ with $\eta \in P$ commute with each other, they preserve each other's generalized eigenspaces, and we can find a simultaneous Jordan decomposition: There is a finite set $\Lambda$ of linear forms $\lambda: P \rightarrow \mathbb{C}$, invariant under complex conjugation $(\lambda \in \Lambda \Rightarrow \bar{\lambda} \in \Lambda$ ), and a decomposition

$$
\begin{equation*}
V_{-}^{c}=\sum_{\lambda \in \Lambda} V_{\lambda} \tag{27}
\end{equation*}
$$

such that $V_{\lambda} \subset \operatorname{ker}\left(S_{\eta}-\lambda(\eta) I\right)^{k}$ for all $\eta \in P$, where $k \in \mathbb{N}$ is sufficiently large. In fact, if $\eta_{1}, \ldots, \eta_{r}$ is a basis of $P$, then $V_{\lambda}=\bigcap_{j=1}^{r} E_{j}$ where $E_{j}$ is the generalized eigenspace of $S_{\eta_{j}}$ corresponding to the eigenvalue $\lambda\left(\eta_{j}\right)$. Since $S_{\eta_{j}}$ is parallel, the same holds for the generalized eigenspaces $E_{i}$ and their intersection, thus the decomposition (27) is parallel along $M$ and therefore $\mathfrak{k}_{+}$-invariant.

[^1]Moreover, each $V_{\lambda}$ is nondegenerate which we see by induction over $r$ : If $r=1$, then $V_{\lambda}$ is a generalized eigenspace of $S_{\eta}$ and hence nondegenerate. For $r>1$ we put $W=\bigcap_{j=1}^{r-1} E_{j}$. This is nondegenerate by induction hypothesis and invariant under $S_{\eta_{r}}$. Then $V_{\lambda}$ is a generalized eigenspace of $\left.S_{\eta_{r}}\right|_{W}$ (considered as a symmetric endomorphism on $W$ ), and thus $V_{\lambda}$ is a nondegenerate subspace of $W$.

The decomposition (27) is invariant under all $S_{\xi}, \xi \in V_{+}$, and if we combine $V_{\lambda}$ and $V_{\bar{\lambda}}$, we get a real decomposition. Hence we can use again Moore's Lemma and the indecomposability of $M$ to conclude that there is just one such pair, $\Lambda=\{\lambda, \bar{\lambda}\}$, or in other words

$$
\begin{equation*}
V_{-}^{c}=V_{\lambda}+V_{\bar{\lambda}} \tag{28}
\end{equation*}
$$

Lemma 6.2. For any $\eta \in P$ such that $\lambda(\eta)=t \in \mathbb{R}$, we have $S_{\eta}=t I+N$ where $N$ is nilpotent with $S_{H} N=0$.

Proof. Since $\lambda(\eta)=\overline{\lambda(\eta)}=t$, we have $S_{\eta}=t I+N$ on $V_{-}^{c}=V_{\lambda}+V_{\bar{\lambda}}$ with $N$ nilpotent. Take any normal vector $\xi=\alpha(v, w)$. Note that $S_{\xi}$ commutes with $S_{\eta}$ and hence with $N$, thus $S_{\xi} N$ is nilpotent $\left(\left(S_{\xi} N\right)^{k}=S_{\xi}^{k} N^{k}=0\right)$ and so it has trace zero. Now from (23) we obtain

$$
\begin{aligned}
B(\xi, \eta) & =-2 \operatorname{tr}\left(S_{\xi} S_{\eta}\right) \\
& =-2 t \operatorname{tr} S_{\xi} \\
& =-2 t \sum \epsilon_{i}\left\langle S_{\xi} e_{i}, e_{i}\right\rangle \\
& =-2 t \sum \epsilon_{i}\left\langle\xi, \alpha\left(e_{i}, e_{i}\right)\right\rangle \\
& =-2 t\langle\xi, H\rangle
\end{aligned}
$$

On the other hand, from (22) we see

$$
\begin{aligned}
B(\xi, \eta) & =-2\left\langle\alpha\left(S_{\eta} v, w\right), H\right\rangle \\
& =-2 t\langle\alpha(v, w), H\rangle-2\langle\alpha(N v, w), H\rangle \\
& =-2 t\langle\xi, H\rangle-2\left\langle S_{H} N v, w\right\rangle .
\end{aligned}
$$

Comparing the two results we see that $S_{H} N=0$
Lemma 6.3. Either $S_{H}^{2}=0$ or there is some $\eta \in P$ with $S_{\eta}=t I$ for some nonzero $t \in \mathbb{R}$.

Proof. Suppose that $S_{H}$ has a non-real eigenvalue $\lambda(H)$. Then $S_{H}^{2}$ and $S_{H}$ are linearly independent, and so $H_{1}$ and $H$ are linearly independent. Then we find a real nonzero linear combination $\eta=a H+b H_{1}$ such that $\lambda(\eta)=t$ is real. From the previous lemma we see that $S_{\eta}=t I+N$ with $S_{H} N=0$. But $S_{H}$ is invertible, hence $N=0$. We have $S_{\eta}=t I$ with $t \neq 0$ : If $S_{\eta}=0$ for any parallel normal field $\eta$, we would have $d \eta=0$ and thus $\eta$ would be a constant vector with $M \subset \eta^{\perp}$, but this was excluded by the fullness assumption.

On the other hand, if $\lambda(H) \in \mathbb{R}$, we have $S_{H}=s I+N$ with $S_{H} N=0$, but $S_{H} N=s N+N^{2}$. If $N^{2} \neq 0$, then $N^{2}$ and $N$ are linearly independent and $N^{2}+s N=0$ is impossible. Thus $N^{2}=0$ and $s=0$.

Now we have finished the proof of Theorem B: $M$ is of Ferus type unless $S_{H}^{2}=0$. In fact, by Lemma 6.3 the affine normal spaces $x+N_{x} M$ have a common intersection point $x-\eta(x) / t$, and the result follows from Theorem 4.1.

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[^1]:    ${ }^{1}$ One way to see this is using the Grassmann valued Gauss map $\tau: M \rightarrow G r_{m}(V)$ which assigns to each $x \in M$ its tangent plane $T_{x} M \subset V$. This is a $K$-equivariant map with $d \tau=\alpha$. Hence ker $d \tau$ is an integrable distribution whose leaves form a euclidean factor, but this cannot hold in the indecomposable case. Thus $\tau$ is an immersion, and from the extrinsic symmetry we see that $\tau(M)$ is a totally geodesic symmetric submanifold of $G r_{m}(V)$. Thus it has nonnegative sectional curvature which is not possible for a minimal submanifold in euclidean space.
    ${ }^{2}$ The integer $k$ is large enough to make $\operatorname{ker} T$ and $\operatorname{im} T$ transversal, and by self adjointness, $\operatorname{ker} T \perp \operatorname{im} T$, hence these subspace form an $S_{H}$-invariant orthogonal decomposition.

